

ELECTROMAGNETISM + MAXWELL'S EQUATIONS

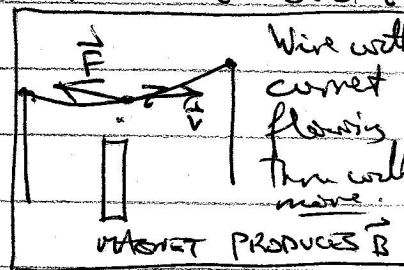
REF FEYNMAN'S LECTURES ON PHYSICS VOL II

Electromagnetism describes how (moving) electric charges interact with each other

EXPERIMENTS SHOW THAT

The force  $\vec{F}$  on a charge  $q$  at location  $\vec{x}$  that has velocity  $\vec{v}$  due to its interaction with other stationary/moving charges is

$$\vec{F} = q (\vec{E} + \vec{v} \times \vec{B}) \quad (1)$$



where  $\vec{E}$ ,  $\vec{B}$  are the Electric + Magnetic (Vector) Fields generated by the other charges (and  $\vec{B}$  depends on  $\vec{v}$  due to all other charges)

NOTICE if  $\vec{v} = \vec{0}$  (ELECTROSTATICS),  $\vec{E} = \vec{F}/q = \text{FORCE PER UNIT CHARGE}$  particle  
 IF  $\vec{x}(t)$  is the position of our charged particle as a function of time  $t$ , then Newton's 2nd Law says

$$m \frac{d^2 \vec{x}}{dt^2} = \vec{F} \quad (2) \text{ where } m \text{ is mass of particle}$$

So if we know  $\vec{E}$ ,  $\vec{B}$  we can ~~solve~~ <sup>use</sup> (1) and (2) to get ODE for motion of our charge

$$m \frac{d^2 \vec{x}}{dt^2} = q \left( \vec{E} + \frac{d\vec{x}}{dt} \times \vec{B} \right)$$

Using ideas from ODEs course we can solve for  $\vec{x}(t)$ .

How do we find  $\vec{E}$  and  $\vec{B}$ ?

- First we need ~~the~~ equations that relate  $\vec{E}, \vec{B}$  to the (other) moving charges. (MAXWELL'S EQNS)
- Then we need to convert Maxwell's equations into equations we can solve and understand (WAVE EQNS)

To understand Maxwell's eqns recall:

① A closed surface  $S$  is one for which  $\partial S = \emptyset$ . Usually  $S = \partial V, \partial E, \partial \Omega$

① The Flux of a V.F.  $\vec{F}$  across a <sup>closed</sup> surface  $S$   $\oint_S$

= Net outward flow across  $S$

$$= \iint_S (\vec{F} \cdot \vec{n}) dS = \iint_S \vec{F} \cdot d\vec{S}$$

② The CIRCULATION of  $\vec{V}$  around a closed loop  $C$

= Net rotational motion around  $C$

$$= \int_C (\vec{F} \cdot \vec{T}) ds = \int_C \vec{F} \cdot d\vec{s}$$

Maxwell's equations is a set of 4 equations. For each one there is a differential equation (DE), or equivalently integral equation (IE) and a physical interpretation (PI)

We describe an expt that motivates the PI

We state PI, Formulate as IE and then show how a version of FTC ~~connects~~ relates the DE to IE.

I GAUSS'S LAW

PI The flux of  $\vec{E}$  through any closed surface  $S$   
=  $\frac{\text{Net Charge inside } S}{\epsilon_0}$   $\epsilon_0 = \text{A constant}$

IE 
$$\iint_S \vec{E} \cdot d\vec{S} = \frac{Q}{\epsilon_0}$$

where  $Q = \text{Net Charge inside } E$ , with  $\partial E = S$

DEF ~~DE~~ Let  $\rho = \rho(x, y, z) = \text{Charge per unit volume}$   
 $= \text{Charge Density.}$

Then

$$Q = \iiint_E \rho \, dV \text{ is net charge in } E.$$

DE The DE form of Gauss' Law is Then

$$\boxed{\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}} \quad \text{Max I}$$

DE  $\Rightarrow$  IE

Suppose  $\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$

Let  $E$  be a solid region with  $\partial E = S$

Then

$$\iiint_E (\nabla \cdot \vec{E}) \, dV = \frac{1}{\epsilon_0} \iiint_E \rho \, dV$$

So by Gauss' Theorem

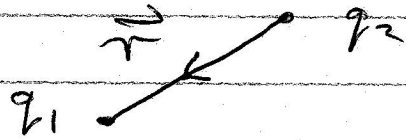
$$\iint_E \vec{E} \cdot d\vec{r}$$

$\perp \cap$

Where does (PI) come from?

Coulomb's Law (EXPT)

The force  $\vec{F}$  on charge  $q_1$  due to charge  $q_2$  is



$$\vec{F} = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2 \vec{r}}{r^3}$$

INVERSE SQUARE LAW.

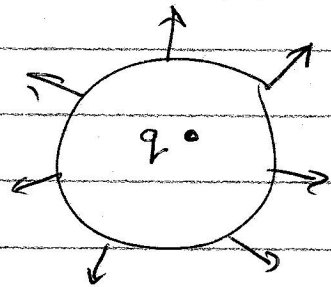
where  $\vec{r}$  = Vector from  $q_2$  to  $q_1$

So Electric Field  $\vec{E}$  due to a charge  $q=q_2$  is  $\vec{E} = \vec{F}/q_1$

$$\vec{E} = \frac{1}{4\pi\epsilon_0} q \frac{\vec{r}}{r^3} \quad (3)$$

- Radial V.F. with length

$$|\vec{E}| = \frac{q}{4\pi\epsilon_0 r^2}$$



PRINCIPLE OF SUPERPOSITION

EF due to a collection of charges is sum of EFs due to each charge

Using This (PI) follows from: let  $\vec{E}$  be EF due to single charge  $q$ . Then

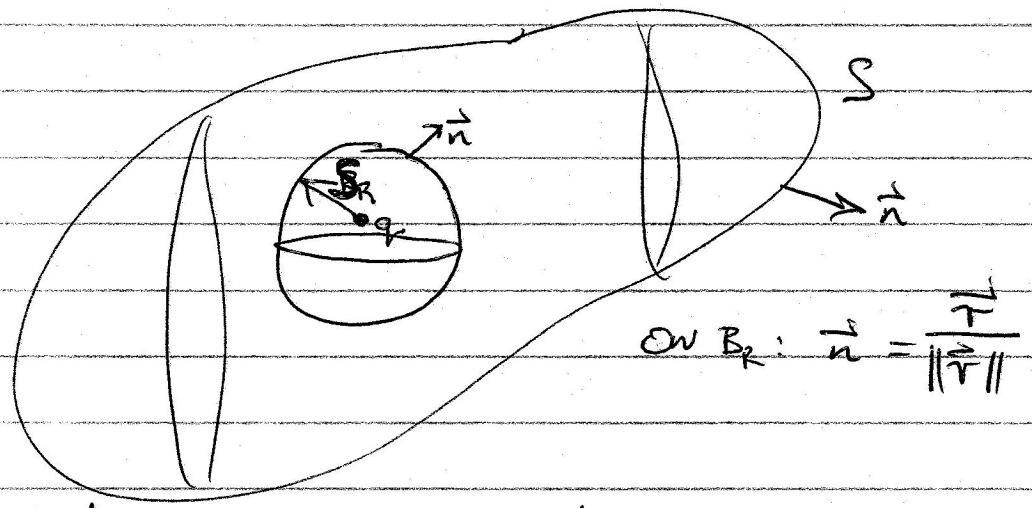
$$\iint_S \vec{E} \cdot d\vec{S} = \begin{cases} 0 & q \text{ outside } S \\ \frac{q}{\epsilon_0} & q \text{ inside } S \end{cases}$$



PROOF OF 4

(A) Put  $q$  at  $(0,0,0)$  inside  $S$

Using  $\vec{r} = (x,y,z)$  you can check  $\nabla \cdot \vec{E} = 0$  on  $\mathbb{R}^3 \setminus \{0\}$ .



Let  $E$  be solid between  $B_R$  and  $S$ ,  $\partial E = S - B_R$   
 Hence  $\nabla \cdot \vec{E} = 0$  on  $E$  we have

$$0 = \iiint_E \nabla \cdot \vec{E} dV \stackrel{\text{GAUSS}}{=} \iint_{\partial E} \vec{E} \cdot d\vec{S} = \iint_S \vec{E} \cdot d\vec{S} - \iint_{B_R} \vec{E} \cdot d\vec{S}$$

So

$$\begin{aligned} \iint_S \vec{E} \cdot d\vec{S} &= \iint_{B_R} \vec{E} \cdot d\vec{S} = \frac{q}{4\pi\epsilon_0} \iint_{B_R} \frac{\vec{r}}{\|\vec{r}\|^3} \cdot \frac{\vec{r}}{\|\vec{r}\|} dS \\ &= \frac{q}{4\pi\epsilon_0} \iint_{B_R} \frac{1}{\|\vec{r}\|^2} dS = \frac{q}{4\pi R^2 \epsilon_0} \iint_{B_R} dS \\ &= \frac{q}{\epsilon_0} \end{aligned}$$

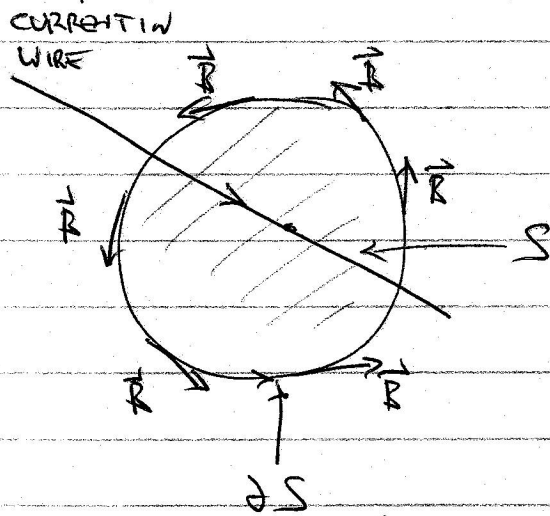
(B) If  $q$  outside  $S$ ,  $\nabla \cdot \vec{E} = 0$  on  $E$  so Gauss' Thm says

$$\iint_S \vec{E} \cdot d\vec{S} = \iiint_E \nabla \cdot \vec{E} dV = 0.$$

II AMPERE'S LAW

A current  $I$  (i.e. moving charges) produces a magnetic field  $\vec{B}$  :

MOVE TO ⊗



Let  $S$  be a surface with boundary  $dS$

(PI) The Circulation of  $\vec{B}$  around  $dS = \frac{\text{Current through } S}{\epsilon_0 c^2}$

where  $c = \text{speed of light}$

LAW OF CONSERVATION OF CHARGE

DEF  $\vec{j} = \text{Current Density}$   
 $\vec{j} / \|\vec{j}\| = \text{DIRECTION of Current}$

$\|\vec{j}\| = \text{Charge per unit area perpendicular to direction of current per unit time.}$

So  $I = \text{Electric Current through } S$   
 $= \text{Total Charge crossing } S \text{ per unit time}$   
 $= \iint_S (\vec{j} \cdot \vec{n}) dS = \iint_S \vec{j} \cdot d\vec{S}$

Conservation of Charge  $\Rightarrow I = - \frac{dQ}{dt}$

$$S_0 \text{ DE-S: } \int_S \vec{j} \cdot d\vec{S} = -\frac{d}{dt} \int_E \rho dV$$

$$\int_E (\nabla \cdot \vec{j}) dV$$

$$S_0 \quad \boxed{\nabla \cdot \vec{j} = -\frac{\partial \rho}{\partial t}} \quad (5)$$

Div of Current Density = -Rate of Charge Density.

⊕

(IE)

$$\int_{\partial S} \vec{B} \cdot d\vec{S} = \frac{1}{\epsilon_0 c^2} \int_S \vec{j} \cdot d\vec{S}$$

Since  $\int_{\partial S} \vec{B} \cdot d\vec{S} \stackrel{\text{STOKES}}{=} \int_S (\nabla \times \vec{B}) \cdot d\vec{S}$  ~~we~~ we get

(DE)

$$\boxed{\nabla \times \vec{B} = \frac{\vec{j}}{\epsilon_0 c^2}} \quad (6)$$

MAXWELL WORCED a Problem (5) + (6)

For any VF  $\vec{R}$ :  $\nabla \cdot (\nabla \times \vec{R}) = 0$ .

S<sub>0</sub>

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot \vec{j} = -\epsilon_0 c^2 \nabla \cdot (\nabla \times \vec{B}) = 0$$

which says CHARGES NEVER MOVE. Whoops!

Maxwell's Correction to (6) is:

$$c^2 \nabla \times \vec{B} = \frac{\vec{J}}{\epsilon_0} + \frac{\partial \vec{E}}{\partial t} \quad \text{Max II}$$

So can produce a Mag Field  $\vec{B}$  from

- Moving charges
- OR • A Time varying electric field.

### (III) No Magnetic Charges Law

There is no analogue of a charge for Mag Fields.

(PI) The flux of  $\vec{B}$  through closed  $S = 0$

$$(IE) \quad \oiint_S \vec{B} \cdot d\vec{S} = 0$$

By Gauss' ~~Theo~~ This gives

$$(PE) \quad \boxed{\nabla \cdot \vec{B} = 0} \quad \text{Max III}$$

### (IV) Faraday's Law

EAT A time varying  $\vec{B}$  produces an EF.

- If move a magnetic around a wire, <sup>loop</sup> induce a current in the wire.



(PI) Circulation of  $\vec{E}$  around loop  $C$

$$= -\frac{d}{dt} \text{Flux of } \vec{B} \text{ Through } S \text{ where } \partial S = C.$$

(IE) 
$$\int_{\partial S} \vec{E} \cdot d\vec{s} = -\frac{d}{dt} \int_S \vec{B} \cdot d\vec{S}$$

By Stokes' Thm 
$$\int_{\partial S} \vec{E} \cdot d\vec{s} = \int_S (\nabla \times \vec{E}) \cdot d\vec{S}$$
  
 so repeat

(DE) 
$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad \text{Max IV}$$

SUMMARY

$$\begin{aligned} \nabla \cdot \vec{E} &= \rho/\epsilon_0 & \textcircled{1} & \qquad \nabla \cdot \vec{B} = 0 & \textcircled{3} \\ \nabla \times \vec{E} &= -\frac{\partial \vec{B}}{\partial t} & \textcircled{2} & \qquad c^2 \nabla \times \vec{B} = \vec{J}/\epsilon_0 + \frac{\partial \vec{E}}{\partial t} & \textcircled{4} \end{aligned}$$

4 PDEs for 2 unknowns  $\vec{E}, \vec{B}$ .

We solve them as follows:

$\nabla \cdot \vec{B} = 0 \Rightarrow \vec{B} = \nabla \times \vec{A}$   $\textcircled{5}$  for some  $\vec{A}$ .

~~Combine with  $\nabla \times \vec{E}$~~

So by  $\textcircled{2}$  
$$\nabla \times \left( \vec{E} + \frac{\partial \vec{A}}{\partial t} \right) = \vec{0}$$
  

$$\vec{E} + \frac{\partial \vec{A}}{\partial t} = -\nabla \phi \quad \text{is conservative}$$

or 
$$\vec{E} = -\frac{\partial \vec{A}}{\partial t} - \nabla \phi \quad \textcircled{6}$$

Using ①, ④ and a clever trick we get eqns for  $\phi, \vec{A}$ :

NOTATION  $\nabla^2 \phi = \nabla \cdot \nabla \phi = \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} + \frac{\partial^2 \phi}{\partial z^2}$

WAVE EQUATIONS

$\nabla^2 \phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = -\frac{\rho}{\epsilon_0}$	⑦
$\nabla^2 \vec{A} - \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2} = -\vec{J} / \epsilon_0 c^2$	⑧

NOTES

① So if know  $\rho, \vec{J}$  can solve wave eqns to get  $\phi, \vec{A}$  and hence  $\vec{E}, \vec{B}$

② ⑦ + ⑧ are examples of a PDE called the WAVE EQUATION. The solution  $\phi(x, y, z, t)$  describes a wave travelling with speed  $c$ .

When  $\rho=0$

One example of sol<sup>n</sup> is  $\phi(x, t) = \sin(x - ct)$

③ Maxwell computed the constant  $c$  using

$$c = \sqrt{\frac{\epsilon_0 c^2}{\epsilon_0}}$$

and found it equaled the speed of light!

This suggested to him that LIGHT <sup>with  $\rho=0, \vec{J}=\vec{0}$</sup>  is a wave ~~field~~ consisting of coupled, time varying  $\vec{E}, \vec{B}$  fields!